TRANSVERSE MOMENTUM DISTRIBUTIONS FROM EFFECTIVE FIELD THEORY

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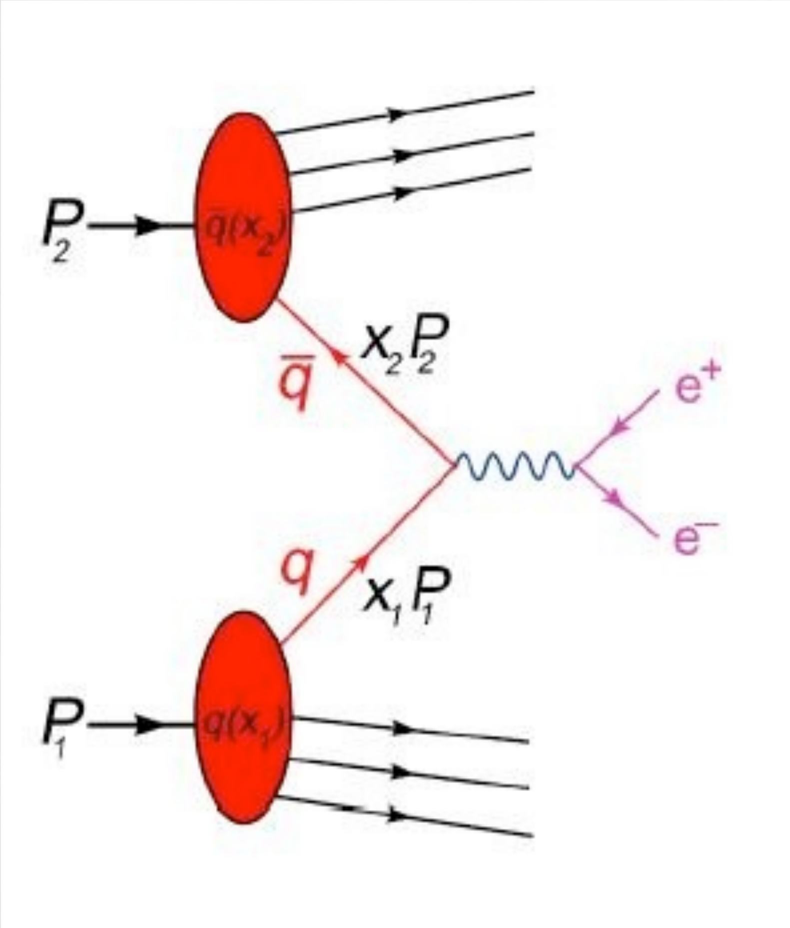
OUTLINE

- Introduction
- The effective field theory approach
- Numerical results and comparison with data
- Conclusion

HADRON COLLIDER

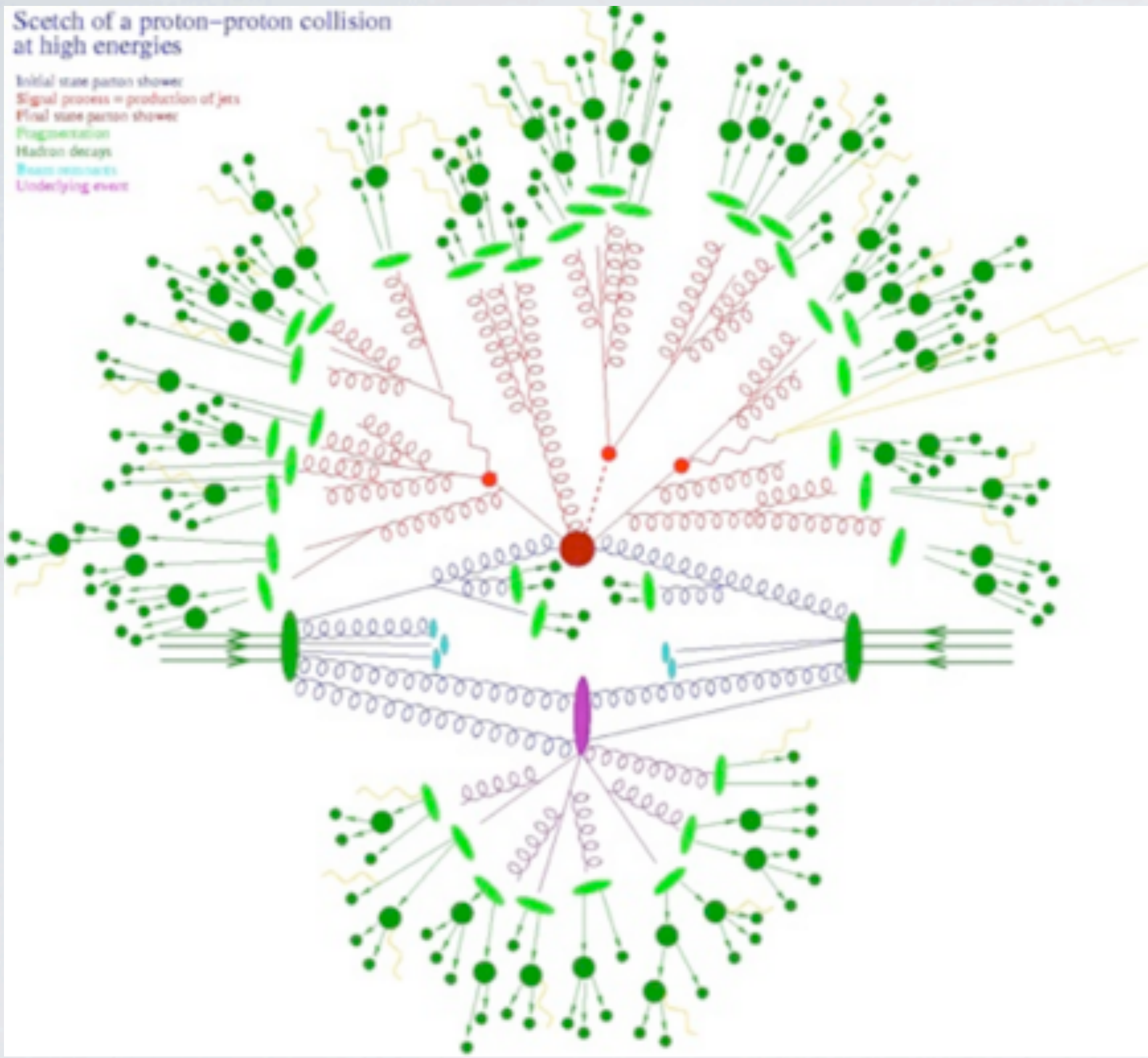
- Useful machine to discover new physics
- Do we really understand what is going on?



?

$$\sigma = \sum_{i,j} \sigma_{i,j}^{part} \otimes f_i(\xi_a) \otimes f_j(\xi_b)$$

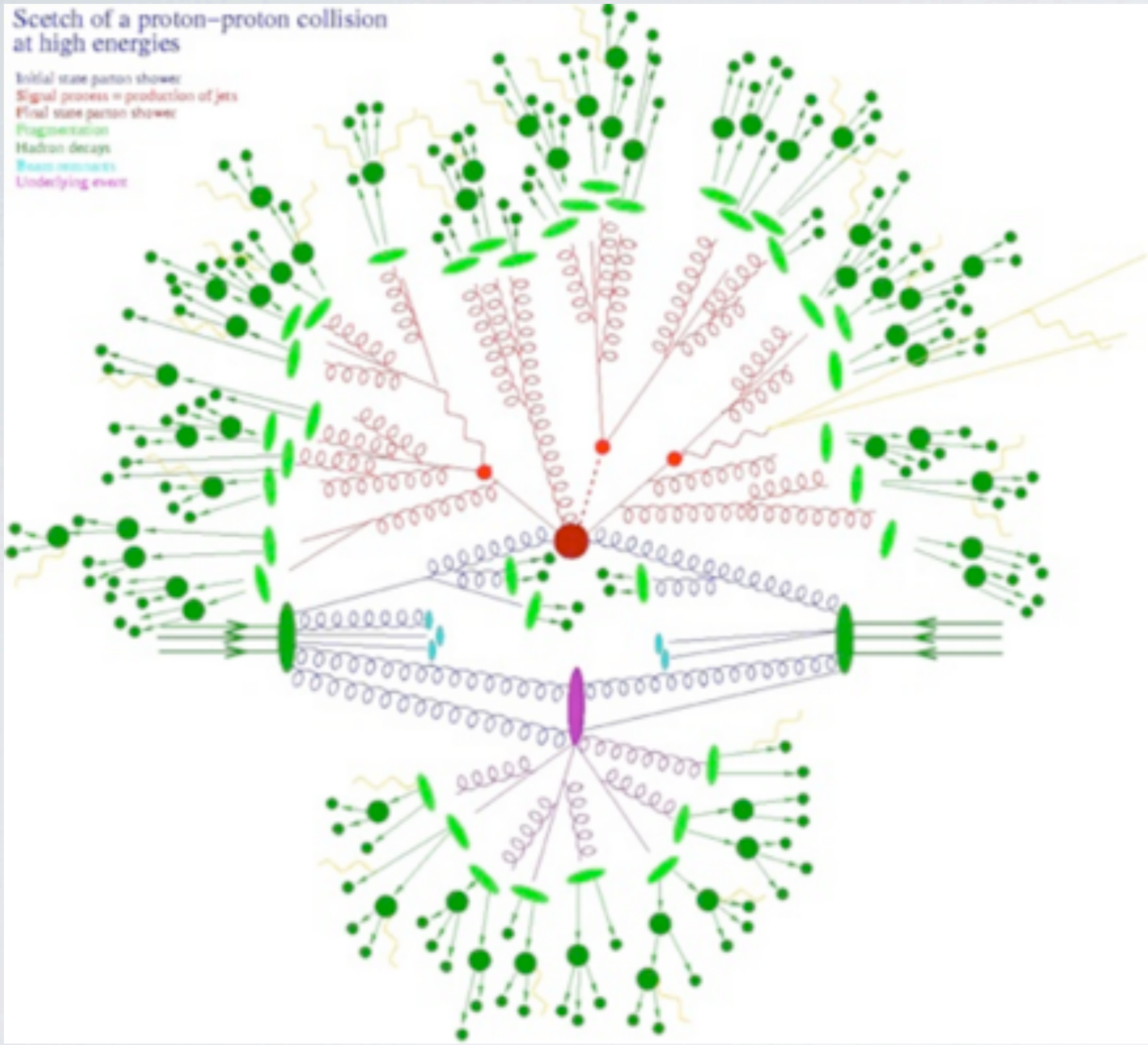
WHAT IS REALLY GOING ON …



- Initial state parton shower
- Signal process (production of jets)
- Final state parton shower
- Fragmentation (hadronization)
- Hadron decays
- Beam remnants
- Underlying events

A Mess !!! Need Factorization

FACTORIZATION



- Physics of interest at hard scale M_H
- Parton shower evolution from M_H to Λ_{QCD}
- Final state hadronization at Λ_{QCD}

Factorization: separates long distance (low energy) and short distance (high energy) behavior

A FAMILIAR EXAMPLE

$$d\sigma = \sum_{i,j} \underbrace{d\sigma_{i,j}^{part}}_{\text{Live at hard scale}} \otimes f_i(\xi_a) \otimes f_j(\xi_b)$$

Live at non-perturbative scale

RG evolve to hard scale

- PDFs live at non-perturbative scale and can be measured experimentally
- Partonic cross section can be obtained using perturbative calculation
- Bring two scales together through RG running

RESUMMATION

$$d\sigma = \sum_{i,j} d\sigma_{i,j}^{part} \otimes f_i(\xi_a) \otimes f_j(\xi_b)$$

Live at hard scale

Live at non-perturbative scale

Kinematic constraints

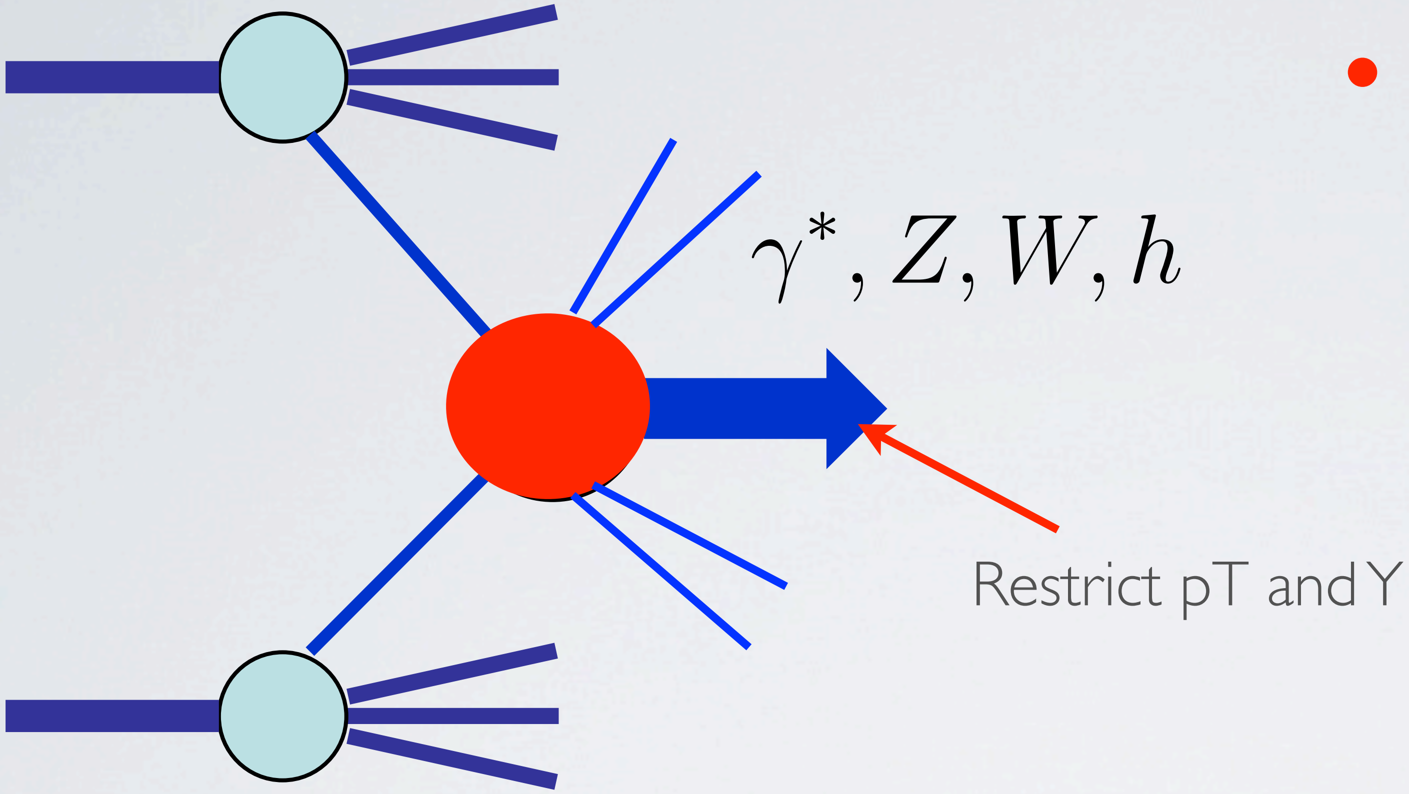
Multiple Disparate scales

RG evolve to hard scale

Additional factorization and resummation required

- Evolution of PDF resums the large logs of hard and non-perturbative scales
- Final state restriction introduces new scales
- Example: low transverse momentum distribution in Drell-Yan process / Higgs production

TRANSVERSE MOMENTUM



- Observable of interest

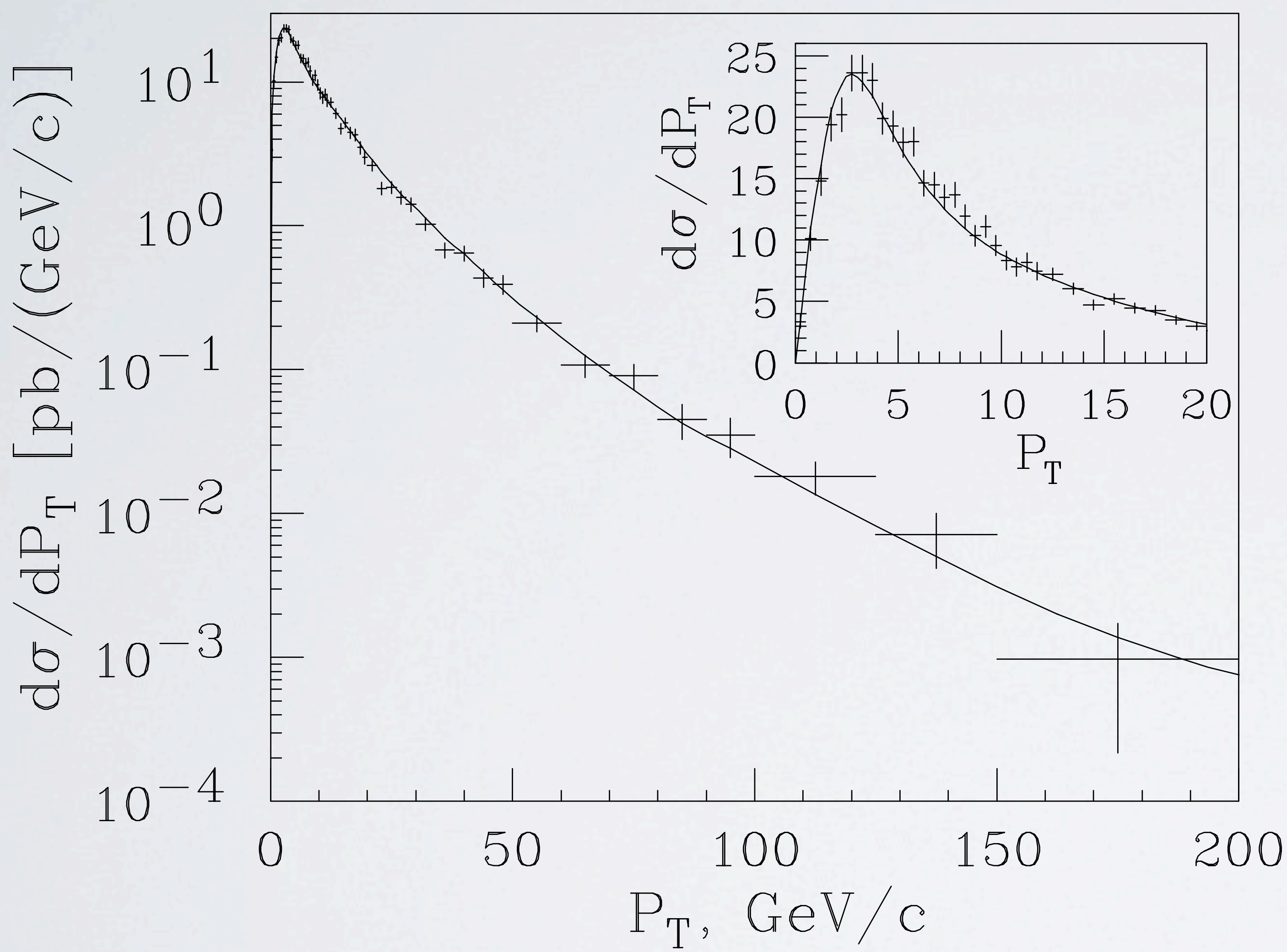
$$\frac{d^2\sigma}{dp_T^2 dY}$$

Motivations

- Higgs Boson searches ➔ pT cut introduced by jet veto
- W-mass measurement ➔ transverse mass endpoint smeared by small W pT due to ISR
- Tests of pQCD
- Probe of transverse nucleon structure

TRANSVERSE MOMENTUM SPECTRUM

Large Logarithms spoil perturbative convergence

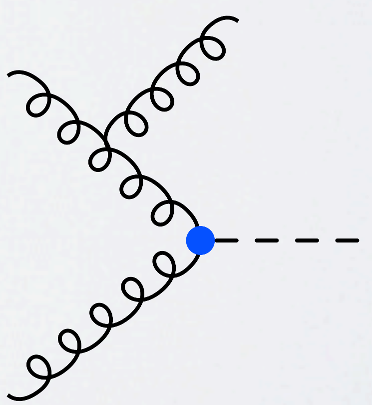


CDF data for Z production:
hep-ex/0001021

- Observable of interest

$$\frac{d\sigma}{dp_T^2} : \frac{1}{p_T^2}\alpha_S^n \ln^k \frac{M_h^2}{p_T^2}$$
$$+\,(\text{non-singular})$$

- High pT region: non-singular term dominates
- Low pT region: perturbation series diverges



RESUMMATION OF TRANSVERSE MOMENTUM

- Resummation has been studied in great detail in the Collins-Soper-Sterman formalism.

(Davies, Stirling; Arnold, Kauffman; Berger, Qiu; Ellis, Veseli, Ross, Webber; Brock, Ladinsky Landry, Nadolsky; Yuan; Fai, Zhang; Catani, Emilio, Trentadue; Hinchliffe, Novae; Florian, Grazzini, Cherdnikov, Stefanis; Belitsky, Ji,....)

- Resummation has also been studied recently using the EFT approach.

(Idilbi, Ji, Juan; Gao, Li, Liu; SM, Petriello; Becher, Neubert)

COLLINS-SOPER-STERMAN FORMALISM

$$\frac{d\sigma_{AB\to CX}}{dQ^2\,dy\,dQ_T^2} = \frac{d\sigma^{(\mathrm{resum})}_{AB\to CX}}{dQ^2\,dy\,dQ_T^2} + \frac{d\sigma^{(\mathrm{Y})}_{AB\to CX}}{dQ^2\,dy\,dQ_T^2}$$

$$\frac{d^2\sigma}{dp_T dY} = \sigma_0 \int \frac{d^2 b_\perp}{(2\pi)^2} e^{-ip_T\cdot b_\perp} \sum_{a,b} [C_a \otimes f_{a/P}](x_A, b_0/b_\perp)[C_a \otimes f_{a/P}](x_B, b_0/b_\perp)$$

$$\times \exp\left\{ \int_{b_0^2/b_\perp^2}^{\hat{Q}^2} \frac{d\mu^2}{\mu^2} \left[\ln\frac{\hat{Q}^2}{\mu^2} A(\alpha_S(\mu^2)) + B(\alpha_S(\mu^2)) \right] \right\}$$

Sudakov Factor

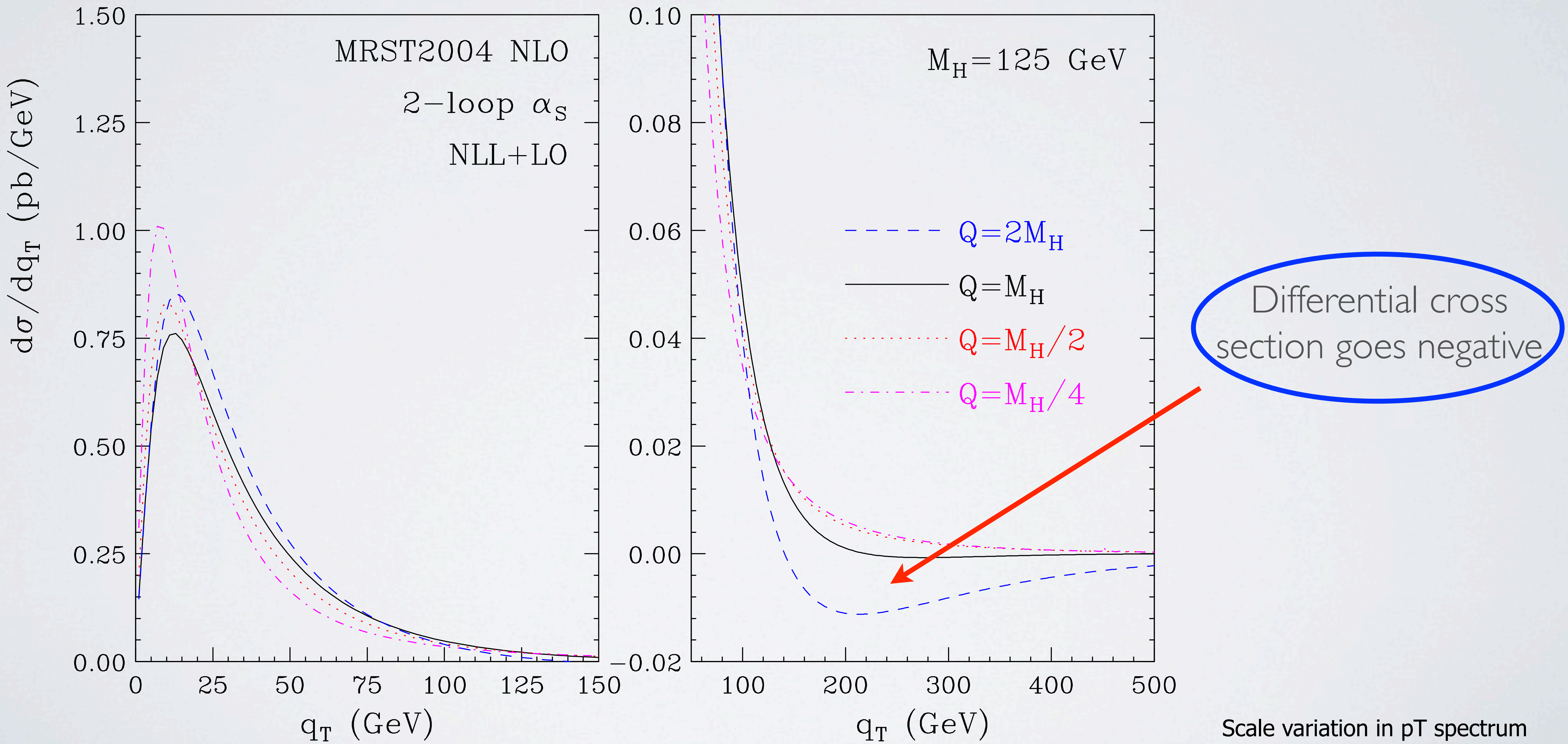
Landau Pole

- Y term neglected for the purpose here
- A,B,C have well-defined perturbative expansions
- Integration of impact parameter b⊥ introduce Landau pole:

a treatment must work for any value of pT

COLLINS-SOPER-STERMAN FORMALISM

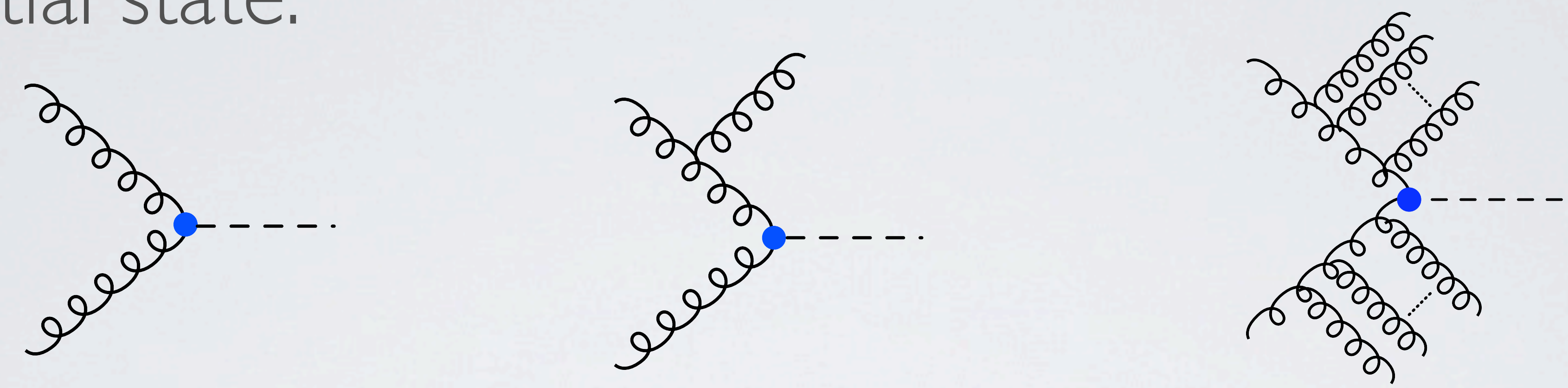
- Resummed exponent in b⊥ space ➔ difficult in matching to fixed order calculation in pT space



Scale variation in pT spectrum using CSS formalism: hep-ph/0508068

EFT FRAMEWORK

- Low pT region dominated by soft and collinear emissions from initial state:

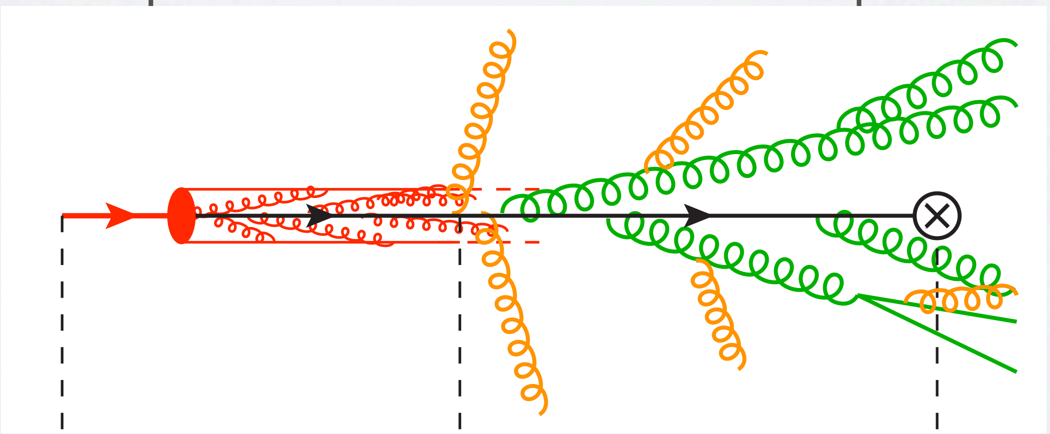


pT=0 finite pT finite pT

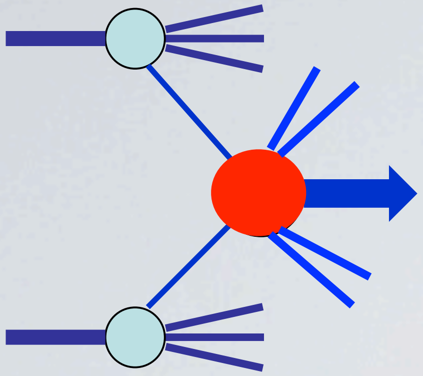
- Hierarchy of scales suggests EFT approach with well defined power counting.

$$m_h \gg p_T \gg \Lambda_{QCD}$$

- Colliding parton is part of initial state pT radiation beam jet:



PDF Beam/Soft Function



EFT FRAMEWORK

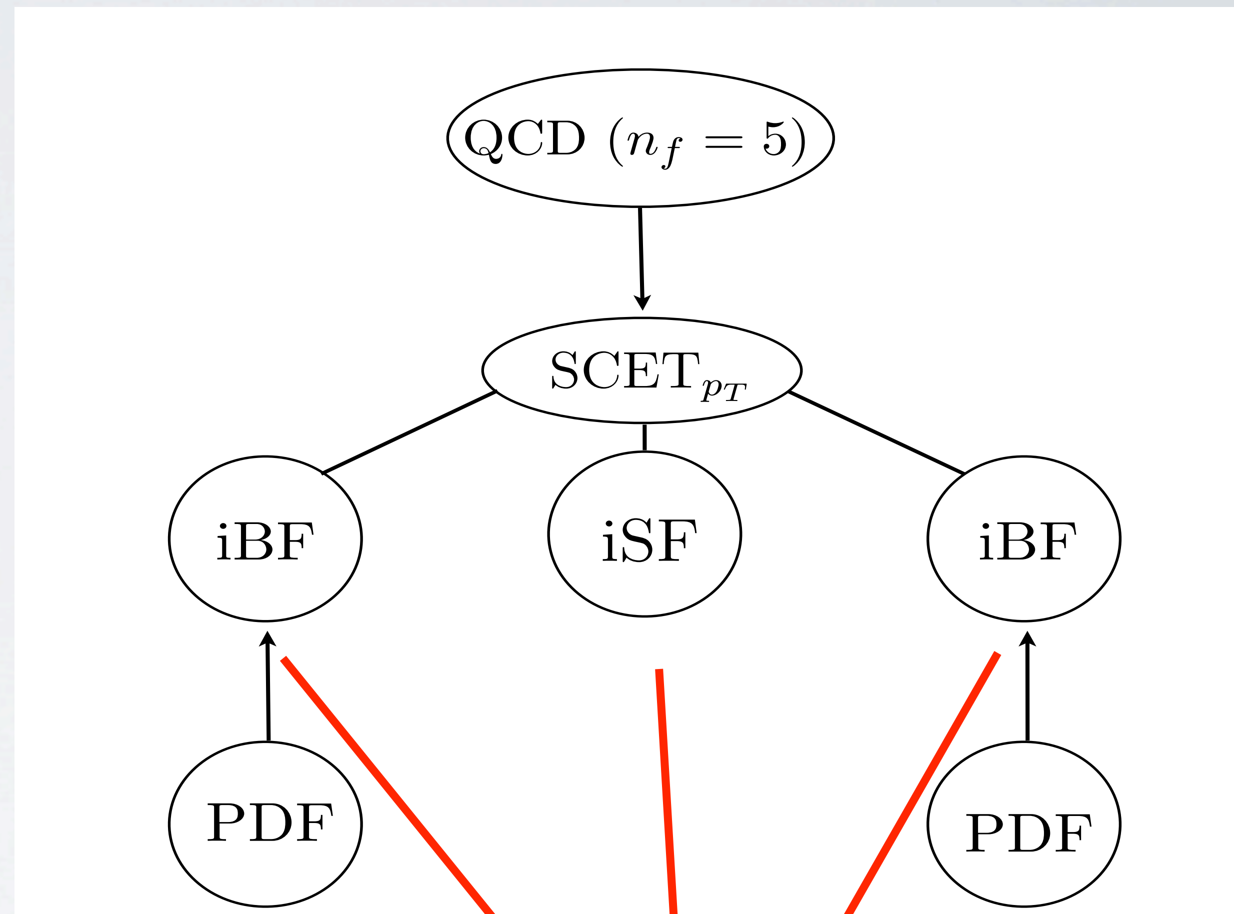
$$\mathrm{QCD}(n_f=6) \to \mathrm{QCD}(n_f=5) \to \mathrm{SCET}_{p_T} \to \mathrm{SCET}_{\Lambda_{QCD}}$$

Top quark integrated out.

Matched onto SCET.

Soft-collinear factorization.

Matching onto PDFs.



Newly defined objects describing soft and collinear pT emissions

$$\frac{d^2\sigma}{dp_T^2\,dY} \sim H \otimes \mathcal{G}^{ij} \otimes f_i \otimes f_j$$

SCET CROSS SECTION

- Schematic form of SCET cross-section:

$$\frac{d^2\sigma}{dp_T^2 dY} \sim \int PS \left| C \otimes \langle \mathcal{O} \rangle \right|^2$$

Wilson coefficient from hard matching

SCET matrix element

- Use soft collinear decoupling to factor out the soft sector

$$\frac{d^2\sigma}{dp_T^2 dY} \sim H \otimes \mathcal{G}^{ij} \otimes f_i \otimes f_j \sim H \otimes B_n \otimes B_{\bar{n}} \otimes S$$

Hard function

Transverse momentum function

Impact-parameter Beam Functions (iBFs): collinear radiation

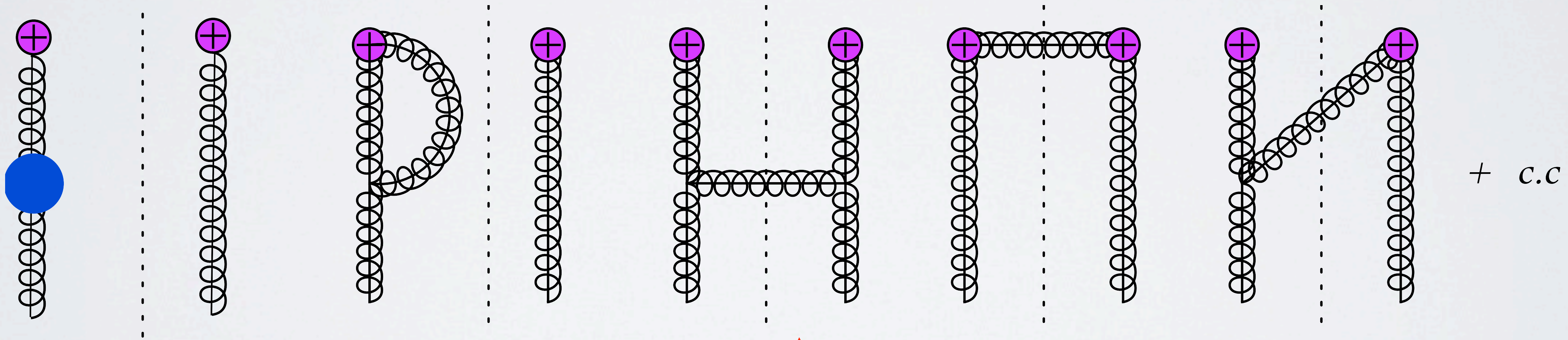
Soft function: soft emission

- Beam function is essentially unintegrated nucleon distribution function and can be matched onto PDF
- The transverse momentum function is a convolution of the iBF matching coefficient and the inverse soft function

IMPACT-PARAMETER BEAM FUNCTION

- Unintegrated nucleon function
- Definition:

$$\tilde{B}_n^{\alpha\beta}(x_1, t_n^+, b_\perp, \mu) = \int \frac{db^-}{4\pi} e^{\frac{i}{2}\frac{t_n^+ b^-}{Q}} \sum_{\text{initial pols.}} \sum_{X_n} \langle p_1 | \big[g B^A_{1n\perp\beta}(b^-, b_\perp) | X_n \rangle \\ \times \langle X_n | \delta(\bar{\mathcal{P}} - x_1 \bar{n}\cdot p_1) g B^A_{1n\perp\alpha}(0) \big] | p_1 \rangle,$$

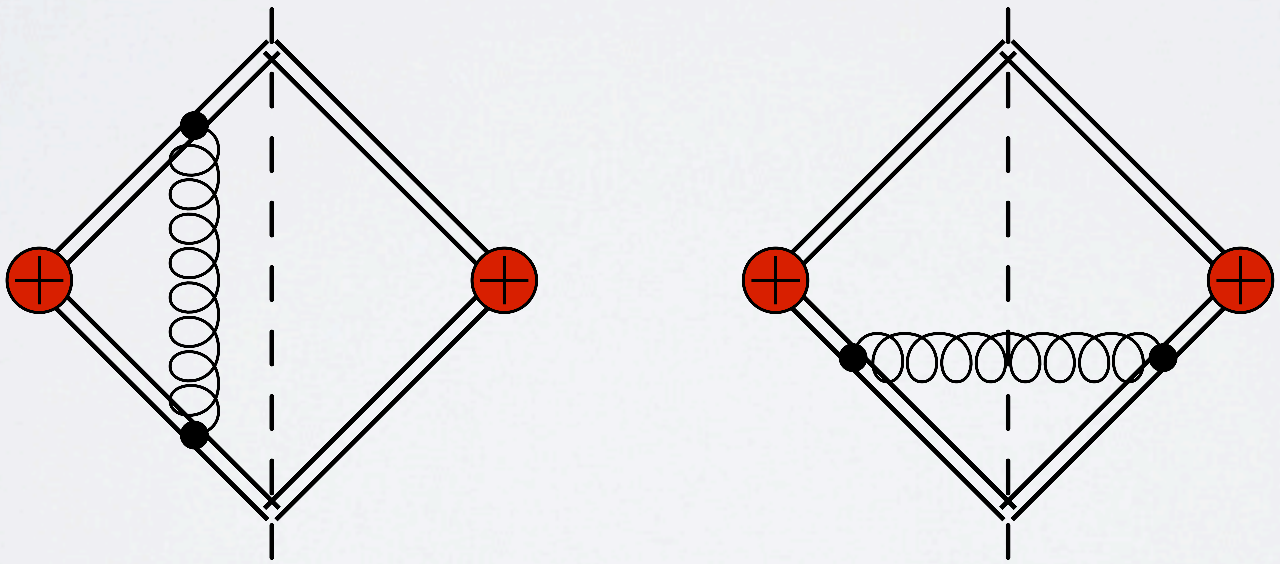


One loop graphs

SOFT FUNCTION

- Unintegrated nucleon function
- Definition:

$$S(z) \;=\; \langle 0|\mathrm{Tr}\big(\bar{T}\{S_{\bar{n}}T^D S^\dagger_{\bar{n}} S_n T^C S^\dagger_n\}\big)(z)\mathrm{Tr}\big(T\{S_n T^C S^\dagger_n S_{\bar{n}} T^D S^\dagger_{\bar{n}}\}\big)(0)|0\rangle$$



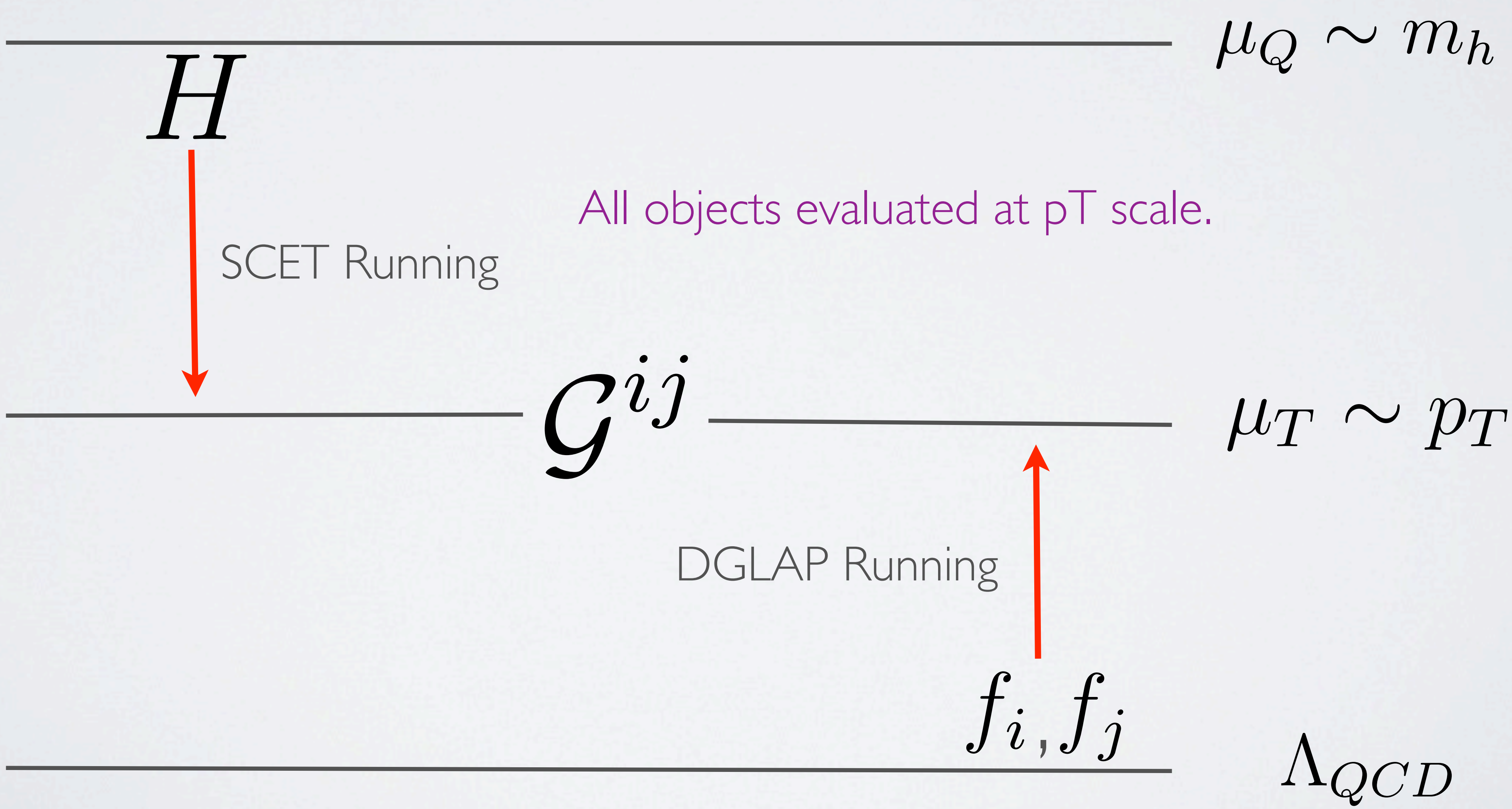
One loop graphs

Running

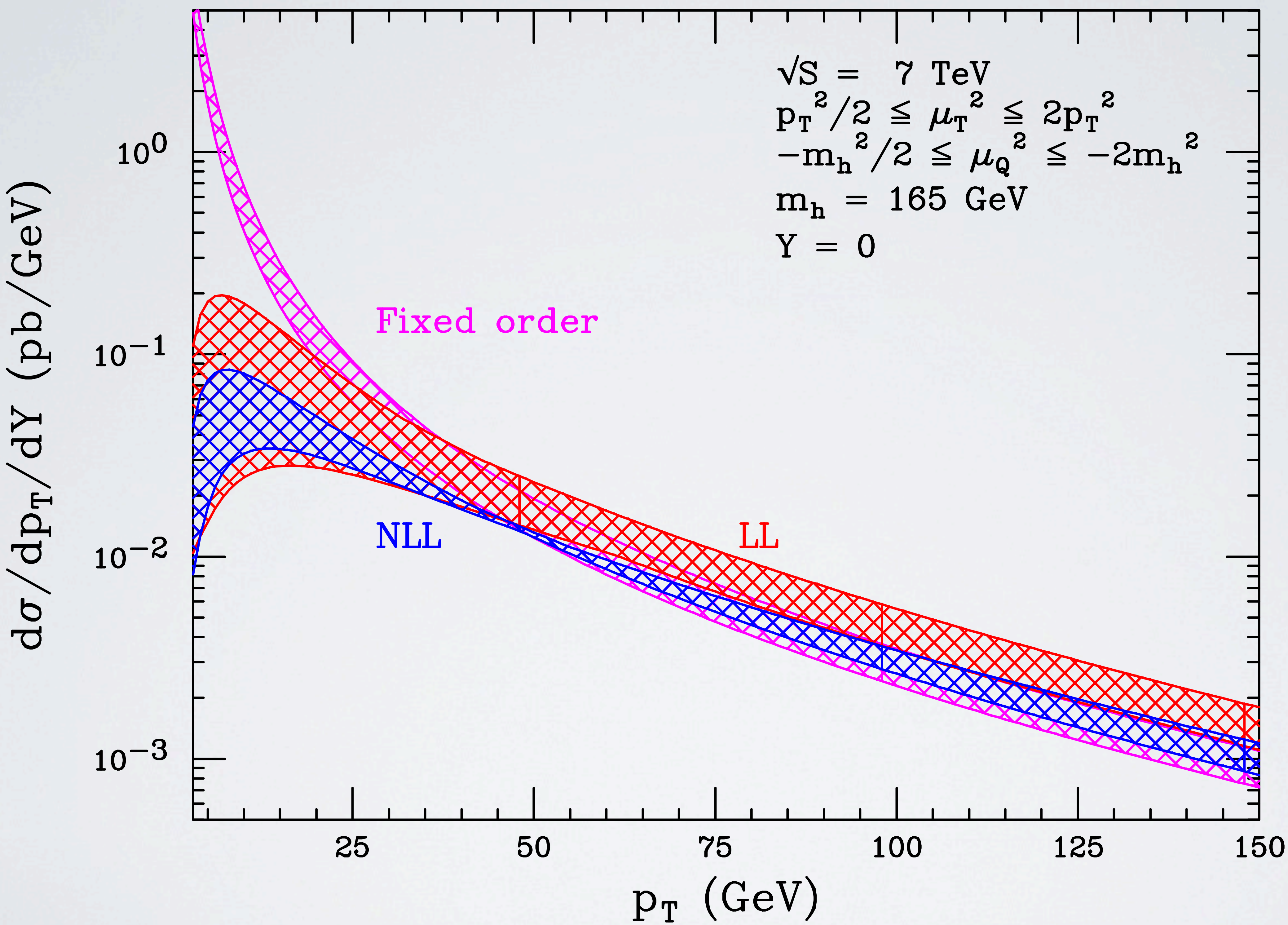
- Factorization formula:

$$\frac{d^2\sigma}{dp_T^2 dY} \sim H \otimes \mathcal{G}^{ij} \otimes f_i \otimes f_j$$

- Schematic picture of running:

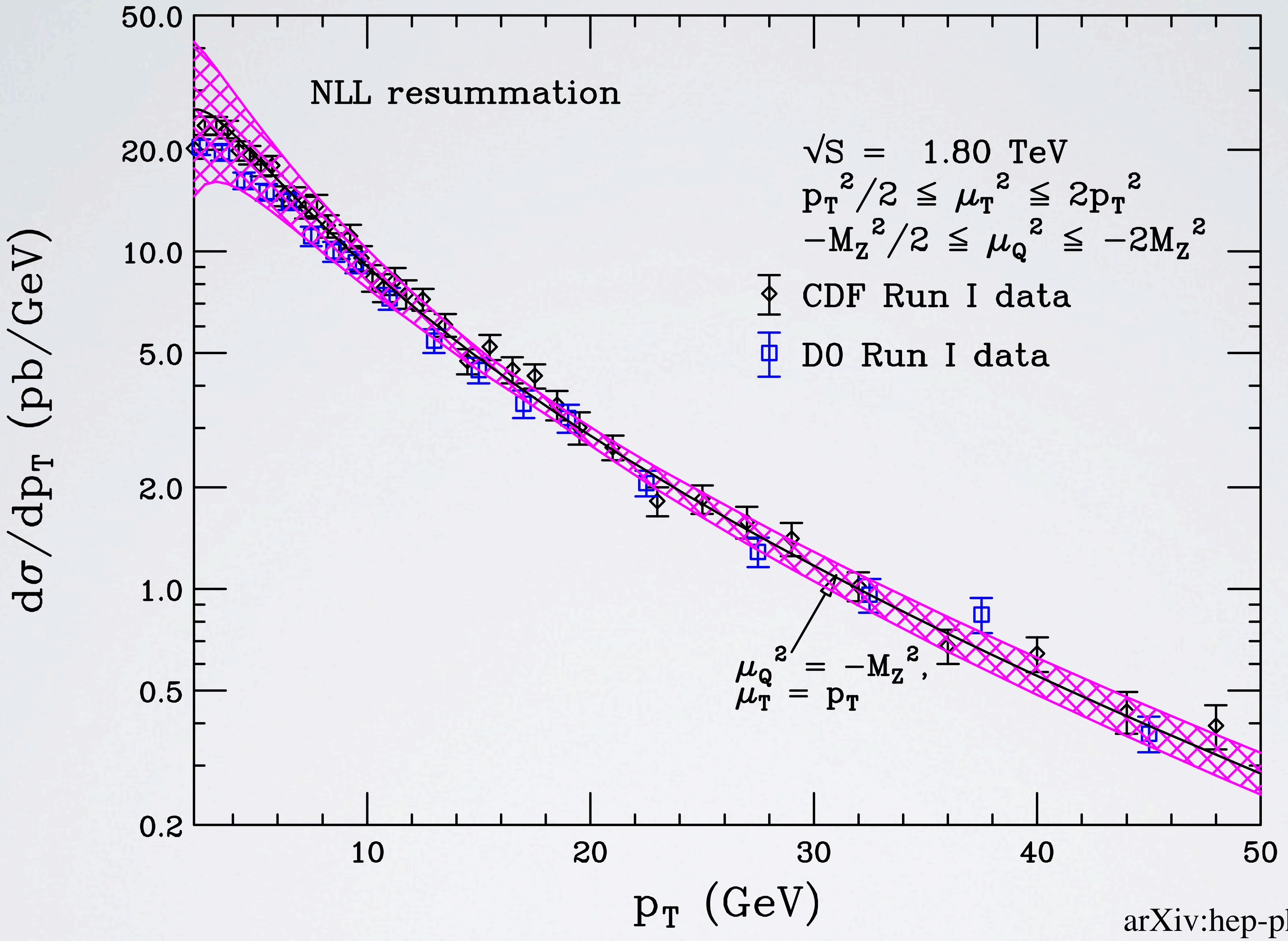


Higgs pT Distribution



- Prediction for Higgs boson pT distribution.

Z-production: Comparison with Data



arXiv:hep-ph/1011.0757

- Good agreement with data.
- Theory curve determined completely by perturbative functions and standard PDFs.

Check to pQCD

- Expand resummed formula to compare to fixed order

$$\frac{d^2\sigma_{Z,q\bar{q}}}{dp_T^2 dY} = \frac{4\pi^2}{3}\frac{\alpha}{\sin^2\theta_W} e_{q\bar{q}}^2 \frac{1}{s\,p_T^2}\sum_{m,n}\left(\frac{\alpha_s(\mu_R)}{2\pi}\right)^n {}_nD_m \ln^m\frac{M_Z^2}{p_T^2}$$

$$\begin{aligned}
\text{leading logarithmic} &: \alpha_s^n L^{2n-1},\\
\text{next-to-leading logarithmic} &: \alpha_s^n L^{2n-2},\\
\text{next-to-next-to-leading logarithmic} &: \alpha_s^n L^{2n-3}.
\end{aligned}$$

Leading Log

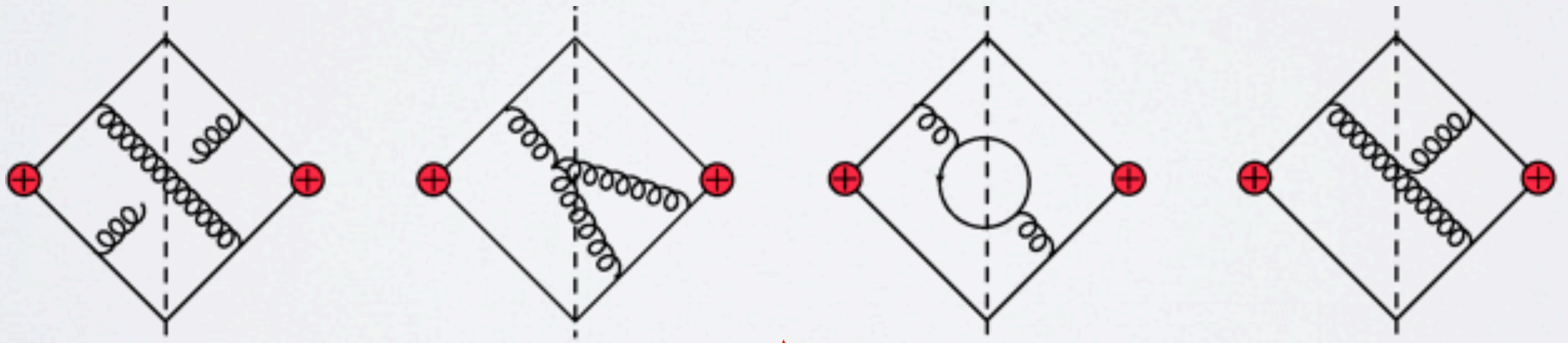
Next-to-Leading Log

$$\begin{aligned}
{}_1D_1 &= A^{(1)} f_A f_B,\\
{}_1D_0 &= B^{(1)} f_A f_B + f_B\left(P_{qq}\otimes f\right)_A + f_A\left(P_{qq}\otimes f\right)_B,\\
{}_2D_3 &= -\frac{1}{2}\left[A^{(1)}\right]^2 f_A f_B,\\
{}_2D_2 &= -\frac{3}{2}A^{(1)}\left[f_B\left(P_{qq}\otimes f\right)_A + f_A\left(P_{qq}\otimes f\right)_B\right] - \left[\frac{3}{2}A^{(1)}B^{(1)} - \beta_0 A^{(1)}\right] f_A f_B,\\
{}_2D_1 &= \left\{-A^{(1)} f_B\left(P_{qq}\otimes f\right)_A \ln\frac{\mu_F^2}{M_Z^2} - 2B^{(1)} f_B\left(P_{qq}\otimes f\right)_A - \frac{1}{2}\left[B^{(1)}\right]^2 f_A f_B\right.\\
&\quad + \frac{\beta_0}{2}A^{(1)} f_A f_B \ln\frac{\mu_R^2}{M_Z^2} + \frac{\beta_0}{2}B^{(1)} f_A f_B - \left(P_{qq}\otimes f\right)_A\left(P_{qq}\otimes f\right)_B\\
&\quad \left. - f_B\left(P_{qq}\otimes P_{qq}\otimes f\right)_A + \beta_0 f_B\left(P_{qq}\otimes f\right)_A\right\} + [A \leftrightarrow B].
\end{aligned}$$

Agrees through NLL level
${}_2D_1$ requires NNLL

Next-to-Next-to Leading Logarithm

- NNLO Beam/Soft function required for NNLL resummation
- Soft function worked out as the first step
- NNLO beam function in progress



arXiv:hep-ph/1105.5171

Two loop graphs for soft function

Soft Function at NNLO

- Anomalous dimensions in position and impact-parameter space
- Old result confirmed: Belitsky (hep-ph/9808389)
- New in impact-parameter space
- New renormalized soft function in full position and impact-parameter space

Define $L = -\frac{b^+ b^- \mu^2 e^{2\gamma_E}}{4}$

$$\gamma_s^{(1)}(b) = 2\frac{\alpha_s}{\pi} C_F \ln(L)$$

$$\gamma_s^{(2)}(b) = \left(\frac{\alpha_s}{\pi}\right)^2 \left\{ C_F N_F \left[-\frac{5}{9}\ln(L) + \frac{\pi^2}{36} - \frac{14}{27}\right] + C_F C_A \left[\left(-\frac{\pi^2}{6} + \frac{67}{18}\right)\ln(L) - \frac{7}{2}\zeta(3) - \frac{11\pi^2}{72} + \frac{101}{27}\right]\right\}$$

Define $L_{0,0} = \delta(q^-)\delta(q^+)$

and $L_{0,1} = \frac{1}{\mu}\left[\frac{\mu}{q^+}\right]_+ \delta(q^-) + \frac{1}{\mu}\left[\frac{\mu}{q^-}\right]_+ \delta(q^+)$

$$\gamma_s^{(1)}(q^-, q^+) = -2\frac{\alpha_s}{\pi} C_F L_{0,1}$$

$$\gamma_s^{(2)}(q^-, q^+) = \left(\frac{\alpha_s}{\pi}\right)^2 \left\{ C_F N_F \left[\frac{5}{9} L_{0,1} + \left(\frac{\pi^2}{36} - \frac{14}{27}\right) L_{0,0}\right] + C_F C_A \left[\left(\frac{\pi^2}{6} - \frac{67}{18}\right) L_{0,1} - \left(\frac{7}{2}\zeta(3) + \frac{11\pi^2}{72} - \frac{101}{27}\right) L_{0,0}\right]\right\}$$

Summary

- Factorization formula:

$$\frac{d^2\sigma}{dp_T^2 dY} \sim H \otimes \mathcal{G}^{ij} \otimes f_i \otimes f_j$$

- Perturbative pT distribution given in terms of perturbatively calculable functions and the standard PDFs.

- Performed NLL resummation and found good agreement with data.

- Next step: NNLL resummation
  - Soft function done
  - Beam function in progress